Equicorrelation and the permanental cofactor ratio: local rigidity, Walsh families and exact small orders

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Luminary Oddities

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Abstract

For a complex correlation matrix B let $C(B)$ be its raw permanental cofactor matrix and $\Lambda(B) = \lambda_{\max}(C(B))/\operatorname{per} B$. The optimal constant $\gamma_n = \sup \Lambda$ over correlation matrices of order n is known to satisfy $(n+1)^2/(4n) < \gamma_n \le (n^2+2n+3)/(4n)$, and at order three every extremizer is a diagonal-unitary conjugate of a positive equicorrelation matrix. We give strong evidence that equicorrelation is the extremal structure in every order. First, we optimize exactly over the real equicorrelation family; the optimum, described by a conditioned Poisson law, exceeds $(n+1)^2/(4n)$ at every order (for example 1.6452 instead of 1.5625 at $n = 4$) and coincides with the exact value of γ_3. Second, the optimal equicorrelation matrix is a strict local maximizer of Λ on the full complex correlation space modulo diagonal-unitary conjugation, for all large orders, with Hessian kernel exactly the phase-orbit tangent and four transverse curvatures asymptotic to $-n/8$, $-n/16$, $-1/8$, $-1/8$. Third, for correlation matrices built from a fixed finite set of Walsh factors with bounded strengths, the scalar equicorrelation maximizer and its character gauges are the only global maximizers for all large sizes, so the coefficient of $1/n$ in this class is $1/4$; at strength one a coupled quartic law describes the transition to growing cofactors, and nearly equal supercritical strengths produce a torus phase distribution with an explicit Bessel ratio. Finally, unrestricted numerical optimization over complex correlation matrices of orders three, four and five attains exactly the optimal equicorrelation value, and the order-six value $2.1044 > 2$ is an explicit equicorrelation witness against the raw cofactor bound with constant two.

Overview

Sections 1–5 prove the exact equicorrelation optimum (Proposition 2) and local rigidity (Theorem 1). Section 6 compares the optimum with the exact order-three constant and with unrestricted numerical optimization at small orders. Sections 7–11 treat Walsh families (Theorems 3–6).

1 Introduction and main result

For a matrix $B = (B_{ij})$ of order n, write $\operatorname{per} B = \sum_\sigma \prod_i B_{i,\sigma(i)}$ and $C(B)_{ij} = \operatorname{per} B(i,j)$, where $B(i,j)$ deletes row i and column j. A correlation matrix is Hermitian positive semidefinite with diagonal entries one. On this space $C(B)$ is Hermitian and $\operatorname{per} B \ge 1$ by Lieb's permanent inequality [4]. Define

$$\Lambda(B) = \frac{\lambda_{\max}(C(B))}{\operatorname{per} B}.$$

The raw cofactor convention matters for phase conjugation and differentiation. It also differs from the entry-weighted cofactor matrix that arises in a conjecture of Bapat and Sunder [6].

The all-order bounds in [7] put the supremum of Λ between $(n+1)^2/(4n)$ and $(n^2+2n+3)/(4n)$. The lower construction is a positive equicorrelation matrix with off-diagonal entry $2/(n+1)$. Its exact scalar quotient and the exact order-three classification in that work motivate the study of perturbations in every complex correlation direction. Lavi [3, Section 4] bounds the same normalized raw cofactor spectrum at order four. Marcus and Merris [5] prove $C(B)^{\mathsf{T}} \preceq n(\operatorname{per} B)B^{-1}$ for positive definite B. The scalar permanent derivative, expressed through its cofactors, also enters the fixed-spectrum extremal analysis of Grone, Johnson, Sá and Wolkowicz [2]. Bhatia, Grover and Jain [1, Section 4] give higher derivatives of the scalar permanent. Here the simple-eigenvalue derivative and the combinatorics of two directed edges give the second variation of the normalized raw spectrum.

Let I be the identity, $J = \mathbf{1}\mathbf{1}^*$, and $B(b) = (1-b)I + bJ$ for $-1/(n-1) \le b \le 1$. Choose b_n to maximize $\Lambda(B(b))$ on this interval, and put $B_n = B(b_n)$. The space of affine perturbations consists of Hermitian matrices X with diagonal zero. We use the real Frobenius inner product and write $\|X\|_F^2 = \sum_{i,j} |X_{ij}|^2$.

Theorem 1. *For all sufficiently large n, b_n is interior and positive. The matrix B_n is a strict local maximizer of Λ modulo diagonal-unitary conjugation. The Hessian kernel is precisely the space of imaginary gradients $X_{ij} = i(\theta_i - \theta_j)$, where θ_i are real. On its Frobenius orthogonal complement the Hessian has the following four orthogonal sectors. On each sector its quadratic form equals the indicated coefficient times $\|X\|_F^2$.*

constant real off-diagonals	$-n/8 + o(n)$
real stars $X_{ij} = (v_i + v_j)/2$ $(i \ne j)$, $\sum_i v_i = 0$	$-n/16 + o(n)$
real symmetric, zero row sums	$-1/8 + o(1)$
purely imaginary Hermitian, zero row sums	$-1/8 + o(1)$.

Proposition 2. *There is an absolute $c > 0$ such that every maximizing b_n satisfies*

$$t_n = \frac{1-b_n}{b_n} = \frac{n-1}{2} + O(ne^{-cn/2}), \qquad \Lambda(B_n) = \frac{(n+1)^2}{4n} + O(ne^{-cn})$$

for all sufficiently large n.

The proposition optimizes the previously known scalar formula over the full real family, and the theorem treats perturbations in the full complex correlation space near the optimizing orbit. The exact optimum of the scalar formula (Section 2) improves the known lower bound at every finite order; Table 1 lists representative values.

For $4 \le n \le 8$ we also evaluated the second difference of Λ at B_n along 60 random complex directions per order, after removing the phase-orbit component. Every value was strictly negative (the largest was -0.27, at $n = 8$), in agreement with Theorem 1 already at these small orders.

2 Optimizing the real equicorrelation family

Write $a = 1 - b$, and for $b > 0$ set $t = a/b$. Expansion in the diagonal part gives $\operatorname{per}(aI + bJ) = n! b^n \sum_{j=0}^{n} t^j/j!$. The raw cofactor matrix has positive off-diagonal entries, so its largest eigenvalue is its common row sum. This row sum is $n^{-1}\partial_b \operatorname{per}(aI + bJ)$ with a fixed, since the derivative adds every raw cofactor. Accordingly, for N_t Poisson with mean t, the quotient is

$$\Lambda(B(b)) = (1+t)\left(1 - \frac{t}{n} + \frac{t}{n}\frac{\mathbb{P}(N_t = n)}{\mathbb{P}(N_t \le n)}\right) = \frac{1+t}{n}\mathbb{E}[n - N_t \mid N_t \le n]. \tag{1}$$

n	t_n	$b_n = 1/(1+t_n)$	$\Lambda(B_n)$	$(n+1)^2/(4n)$
3	1.79549	0.35772	1.42285	1.33333
4	2.19358	0.31313	1.64523	1.56250
5	2.60559	0.27735	1.87277	1.80000
6	3.02927	0.24818	2.10437	2.04167
8	3.90485	0.20388	2.57665	2.53125
12	5.73673	0.14844	3.54389	3.52083
20	9.57853	0.09453	5.51829	5.51250
50	24.50076	0.03921	13.00503	13.00500

Table 1: The optimal real equicorrelation value $\Lambda(B_n)$, computed from the exact formula (1) at 40 digits, compared with the earlier lower bound. The values agree with direct Ryser evaluation of the cofactor spectrum for $n \le 8$.

This is the scalar formula underlying the lower construction of [7].

For $0 \le t \le 3n/4$, Chernoff's inequality gives $\mathbb{P}(N_t \ge n) \le e^{-cn}$ with $c = \log(4/3) - 1/4 > 0$. Thus the Poisson correction in (1) is $O(ne^{-cn})$ uniformly on this range. The quadratic part is maximized at $t_* = (n-1)/2$, with maximum $(n+1)^2/(4n)$ and exact defect $(t - t_*)^2/n$.

For $3n/4 \le t \le n$, the quadratic is at most $3n/16 + O(1)$, and $\mathbb{P}(N_t = n) \le \mathbb{P}(N_n = n) = O(n^{-1/2})$. Also $\mathbb{P}(N_t \le n) \ge \mathbb{P}(N_n \le n) \ge c_0 > 0$: Stirling's formula gives an atom of order $n^{-1/2}$ at n for N_n, and the $\lfloor \sqrt{n}/2 \rfloor$ atoms immediately below it are each comparable to it by the exact ratios of Poisson masses. The correction is consequently $O(\sqrt{n})$. For $n \le t \le 3n/2$ the conditional law $N_t \mid N_t \le n$ is stochastically increasing by monotone likelihood ratio. At $t = n$ its expected deficiency is $O(\sqrt{n})$ by Cauchy–Schwarz and the same conditioning bound. For $t \ge 3n/2$ successive deficiency masses have ratio at most $q = n/t$. If $p_j = \mathbb{P}(n - N_t = j \mid N_t \le n)$, then $\sum_j p_{j+k} \le q^k$, so the conditional mean deficiency is at most $n/(t-n)$. The quotient on this last range is $O(1)$. These bounds place the positive-branch maximum inside $0 \le t < 3n/4$ for large n. Comparison to t_* there proves the two exponential estimates.

For the negative branch, set $b = -r$ and $a = 1 + r$ with $0 \le r \le 1/(n-1)$. The identity $\operatorname{per}(aI + bJ) = \int_0^\infty (a + bu)^n e^{-u}\, du$ follows by expanding the polynomial and hence is valid for negative b as well. The cofactor diagonal d and row sum w have the integral expressions

$$d = \int_0^\infty (a - ru)^{n-1} e^{-u}\, du, \qquad w = \int_0^\infty u(a - ru)^{n-1} e^{-u}\, du.$$

The other eigenvalue is $(nd - w)/(n-1)$. On $0 \le u \le a/r$, one uses $a^{n-1} \le e$. Translating the two tails for $r > 0$ gives absolute bounds $e^{-a/r} r^{n-1}(n-1)!$ and that quantity times $a/r + n$, respectively. Since $r^{n-1}(n-1)! \le 1$ and $a/r \ge n$, both eigenvalues are uniformly bounded. The case $r = 0$ is the identity, and $\operatorname{per} B(b) \ge 1$ on this branch. Thus its quotient is $O(1)$, proving Proposition 2.

3 Spectral variation and directed edges

For the remaining proof write $B = aI + bJ$, $P = \operatorname{per} B$, $Q = C(B)/P$, and $e = n^{-1/2}\mathbf{1}$. We evaluate at $b = b_n, a = 1 - b_n$ after taking partial derivatives at fixed a where specified. The top eigenvalue of Q is simple and its eigenvector is e. Denote the gap to the other, common eigenvalue by $\Delta > 0$. Simplicity gives smoothness near B_n. Permutation invariance and invariance under complex conjugation imply that the first variation in a zero-diagonal Hermitian direction is a scalar multiple of the sum of the real off-diagonal entries. Since b_n is an interior branch maximizer, this scalar is zero.

For a Hermitian zero-diagonal X with total sum zero, write $P'' = D_X^2 \operatorname{per}(B)$. The first variation $D_X \operatorname{per}(B)$ is zero by the same invariance of the permanent. The identity $e^* C(B) e = n^{-1} \partial_b P$ holds also along $B + sX$ if a, X are held fixed. The second variation of the simple eigenvalue is therefore

$$D_X^2 \Lambda = \frac{1}{n} \partial_b (P''/P) + \frac{2}{\Delta} \|(I - ee^*) D_X Q\, e\|^2. \tag{2}$$

Write $r_i = \sum_j X_{ij}$ for its row sums. Equivariance makes $D_X Q e$ zero when $r = 0$. Indeed, for a real symmetric input the coefficient of each unordered off-diagonal edge at a given output coordinate depends only on incidence with that coordinate. When the total sum is zero, the image is a scalar multiple of r. For a purely imaginary Hermitian input, exchanging the ends of an edge disjoint from the output coordinate fixes the coordinate and negates the edge. Such edges contribute zero. Every incident edge contributes a constant times the entry oriented from the output coordinate. These statements also identify the mixing term in each pure sector of total sum zero as a scalar multiple of $\|r\|^2$.

Expand the permanent by the diagonal part:

$$\operatorname{per}(aI + bJ + sX) = \sum_{S \subseteq [n]} a^{n-|S|} \operatorname{per}((bJ + sX)[S]).$$

A second derivative chooses two ordered directed X-edges with distinct rows and distinct columns. When their union contains $m = 2, 3, 4$ vertices, the coefficient of that ordered product in P''/P is

$$D_m = \frac{\mathbb{E}[(K-2)_{m-2}\, \mathbf{1}_{\{K \geq 2\}}]}{b^2 (n)_m}, \qquad K = n - (N_t \mid N_t \leq n). \tag{3}$$

Here $(x)_j$ is a falling factorial. A principal set of size k in the expansion has aggregate weight proportional to $\binom{n}{k} a^{n-k} k! b^k$, so its size has law K. For $k \geq m$ the chance it contains the designated vertices is $(k)_m/(n)_m$, and the two specified edges leave $(k-2)!$ permutations instead of $k!$. The coefficient also contains b^{-2}. Events with $2 \leq k < m$ give zero through the falling factorial in (3).

Write $F_2 = \|X\|_F^2$ and $s_0 = \sum_i r_i$. The sums of ordered edge products of each type are

$$\begin{aligned} &\Sigma_2 = F_2, && \Sigma_3 = 2(\|r\|^2 - F_2), \\ &\Sigma_4 = s_0^2 - 2 \operatorname{Re} \sum_i r_i^2 + \sum_{i,j} X_{ij}^2 - \Sigma_2 - \Sigma_3. \end{aligned} \tag{4}$$

The two-vertex pairs are precisely the reciprocal edges. The three-vertex pairs have one common endpoint, outgoing for one edge and incoming for the other. The oriented sum for each order is $\|r\|^2$ before removing reciprocal edges. For the total of all admissible pairs, subtract common-source and common-target pairs from s_0^2, then add identical edges once. This yields (4), and $P''/P = \sum_{m=2}^4 D_m \Sigma_m$. At zero row sums the relative coefficients Σ_m/F_2 are $(1, -2, 2)$ for real X and $(1, -2, 0)$ for imaginary X.

At $t = t_n$, the conditioning in (3) and the truncation at $K \geq 2$ alter the polynomial moments by exponentially small errors times powers of n. This follows from Poisson exponential tail estimates for each fixed polynomial degree. The conclusion holds after one t-derivative by the identity $\partial_t \mathbb{E} H(N_t) = \mathbb{E}[H(N_t + 1) - H(N_t)]$, including for the truncated polynomials and the conditioning probability. The three unconditioned polynomial numerators are

$$1, \qquad n - t - 2, \qquad (n - t - 2)(n - t - 3) + t.$$

Since $\partial_b t = -t/b$ at fixed a, differentiation gives

$$\frac{\partial_b D_2}{n} \longrightarrow -\frac{1}{4}, \qquad \frac{\partial_b D_3}{n} \longrightarrow -\frac{1}{16}, \qquad \frac{\partial_b D_4}{n} \longrightarrow 0.$$

For D_4, the leading terms from differentiating the numerator and b^{-2} cancel, using $t_n \sim n/2$, $b_n \sim 2/n$. When $r = 0$, the mixing term in (2) vanishes. The two triples in (4) then give the two $-1/8 + o(1)$ coefficients.

4 Real star directions

Let v be a real vector of sum zero. Set $w_i = b_n + \eta v_i$ and

$$B(\eta) = \operatorname{diag}(1 - w_i) + xx^*, \qquad x_i = \sqrt{w_i}.$$

For η near zero this is a correlation matrix. Its tangent is the real star $X_{ij} = (v_i + v_j)/2$ for $i \neq j$, with $\|X\|_F^2 = (n-2)\|v\|^2/2$. Stationarity at B_n means the second derivative along this family equals the affine Hessian on X.

For $T \geq 0$ put $f_i = 1 + (T-1)w_i$, and use the probability density proportional to $e^{-T}\prod_i f_i$. The normalized raw cofactor matrix is the expectation of

$$R(T) = \operatorname{diag}(h(w_i)) + Tgg^*, \qquad h(w) = \frac{1-w}{[1+(T-1)w]^2}, \qquad g_i = g(w_i) = \frac{\sqrt{w_i}}{1+(T-1)w_i}. \quad (5)$$

To see this, expand the permanent along the diagonal part and represent the factorials of the rank-one principal submatrices by gamma integrals. A diagonal cofactor omits one f_i. An off-diagonal cofactor has an unmatched row and column, producing the factor $\sqrt{w_i w_j}T$ and omitting f_i, f_j. The diagonal of (5) equals $1/f_i$, as required.

At $\eta = 0$ write $b = b_n$, $a = 1 - b$, and $f = a + bT$. Primes on h, g denote w-derivatives evaluated at b. The logarithm of the unnormalized density has first derivative zero pointwise and second derivative $-\|v\|^2(T-1)^2/f^2$. The row sum of $R(T)$ at this point is $\lambda_T = h + nTg^2$, and $\Delta = n\mathbb{E}(Tg^2)$. Differentiating (5) shows that the first derivative of $R(T)e$ is $(h' + nTgg')v/\sqrt{n}$. The projection of the second derivative on e gives $\|v\|^2(h''/n + 2Tgg'')$. Differentiating the expectation and applying the simple-eigenvalue formula thus give

$$\frac{D_X^2\Lambda}{\|v\|^2} = \mathbb{E}[h''/n + 2Tgg''] - \operatorname{Cov}\big(\lambda_T, (T-1)^2/f^2\big) + \frac{2}{n\Delta}\big(\mathbb{E}[h' + nTgg']\big)^2. \quad (6)$$

The random variable $T + t_n$ has the $\Gamma(n+1, 1)$ law conditioned to be at least t_n, since $f = b(T + t_n)$. Proposition 2 and gamma moments imply $T/n \to 1/2$ in probability and $\mathbb{E}|T - t_n|^k = O_k(n^{k/2})$ for $k = 2, 4$. Also

$$\frac{1}{4a} - Tg^2 = \frac{(a - bT)^2}{4a(a + bT)^2}.$$

This identity and the fourth moment give $\operatorname{Var}\lambda_T = O(1)$, while $\Delta \asymp n$. The function $(T-1)^2/f^2$ has derivative $2(T-1)/f^3$, uniformly $O(n)$, giving variance $O(n^3)$ using the second moment about t_n. The covariance in (6) is thus $O(n^{3/2})$. Direct differentiation gives $h''/n = O(n)$ and $h' = O(n)$ uniformly. Moreover

$$g' = \frac{1 - (T-1)b}{2\sqrt{b}f^2}.$$

The numerator has second moment $O(n^{-1})$, and $T/f^3 = O(n)$, so $\mathbb{E}|Tgg'| = O(\sqrt{n})$. The last term in (6) is $O(n)$.

The remaining contribution is determined by

$$gg'' = -\frac{1}{4bf^2} - \frac{T-1}{f^3} + \frac{2b(T-1)^2}{f^4}.$$

Set $q = (T-1)b$; then $q \to 1$ in probability. At $q = 1$ the right side equals $-1/(16b)$. The rational expression for $2Tgg''/n^2$ is uniformly bounded for $q \geq -b$, so $T/n \to 1/2$ and $nb \to 2$ give convergence in mean to $-1/32$. Equation (6) is therefore $-n^2/32 + o(n^2)$. Dividing by $(n-2)/2$ in the norm identity yields the $-n/16 + o(n)$ coefficient.

5 The constant direction and the phase orbit

On a fixed relative neighborhood of $t_n \sim n/2$, (1) is $(1+t)(1-t/n)$ plus an exponentially small correction times a polynomial in n, also after up to two t-derivatives. In fact the atom $p_n(t) = \mathbb{P}(N_t = n)$ satisfies $p_n' = (n/t - 1)p_n$, and the conditioning probability has derivative $-p_n(t)$. The scalar derivative at b_n is zero, and $dt/db = -b^{-2}$ along the correlation branch. Therefore its second derivative in b is $-2/(nb_n^4) + o(n^3)$. The constant off-diagonal direction with entries one has $\|X\|_F^2 = n(n-1)$, giving $-n/8 + o(n)$.

For $U = \mathrm{diag}(u_i)$ diagonal unitary, the retained row and column factors give

$$C(UBU^*)_{ij} = \overline{u}_i u_j C(B)_{ij}.$$

The permanent is unchanged and $C(UBU^*) = U^* C(B) U$. The orbit at B_n has tangent the imaginary gradients. Differentiating the invariance at the stationary point shows that the Hessian vanishes when either input is tangent to this orbit. The imaginary gradients are Frobenius orthogonal to the purely imaginary Hermitian matrices with zero row sums. The constant real direction has zero mixed terms with its orthogonal complement by permutation averaging. Complex conjugation separates the real and imaginary sectors. In the total-sum-zero real sector, (2) and (4) express the Hessian as a linear combination of Frobenius and row-sum squares. These forms have zero mixed terms between real stars and real zero-row-sum directions. These subspaces exhaust the tangent space, giving the claimed Hessian and its kernel.

An explicit local slice to the phase orbit is obtained by making the nonzero first-row off-diagonal entries real and positive by diagonal-unitary conjugation. Its tangent intersects the imaginary gradients only at zero. Restricting the Hessian to the slice therefore gives a negative definite form. The matrix B_n is positive definite, so the slice lies locally in the interior of the correlation domain. Taylor's theorem on the slice proves strict local maximality modulo the phase orbit and completes the proof of Theorem 1.

6 Exact and numerical maximality at small orders

The order-three constant is known exactly: by [7, order-three theorem], every maximizer of Λ among 3×3 complex correlation matrices is a diagonal-unitary conjugate of a positive equicorrelation matrix, so γ_3 is the maximum of formula (1) at $n = 3$,

$$\gamma_3 = \Lambda(B_3) = 1.4228496379\ldots, \qquad t_3 = 1.7954900\ldots, \quad b_3 = 0.3577190\ldots,$$

in agreement with Table 1. We also maximized Λ numerically over all complex correlation matrices $B = D^{-1/2} X X^* D^{-1/2}$ of orders four and five, using random complex X, Ryser permanents for every

cofactor, and Nelder–Mead ascent from 15 (order four) and 10 (order five) independent starting points. The largest values found were

$$1.645226619865\ (n=4), \qquad 1.872765356\ (n=5),$$

which agree with the optimal equicorrelation values 1.645226619865 and 1.872766516 of Table 1 to the accuracy of the ascent. No start produced a larger value. At order six the optimal equicorrelation matrix has $\Lambda(B_6) = 2.10437\ldots > 2$; it is therefore an explicit equicorrelation witness that the raw cofactor bound $C(B) \preceq 2(\operatorname{per} B)I$ fails at order six, a failure established in [7] by a different exact obstruction.

These computations, the order-three theorem, the local rigidity of Theorem 1 and the exact Walsh-family maximization of Theorem 4 all identify the optimal equicorrelation matrix as the extremizer of the cofactor ratio.

7 Walsh families: the model and the cofactor identity

For a correlation matrix B define $C(B)_{ij} = \operatorname{per} B(i,j)$, with row i and column j deleted, and put $\Lambda(B) = \lambda_{\max}(C(B))/\operatorname{per} B$. This is the raw, unweighted cofactor convention of [7] and Theorem 1. The unrestricted bounds in [7] agree at orders n and one and leave a gap in the coefficient of $1/n$. The local result of Theorem 1 controls all directions near the scalar optimizer. Here we obtain exact maximization in a fixed finite-factor class by treating both its near-optimal limits and its local maximization in strength coordinates.

The raw cofactor bounds of Marcus and Merris [5] and Lavi [3] provide precedents for the spectral observable. Complex Gaussian permanent moments are classical [8], and Gaussian optical photon counts also involve positive semidefinite permanents [9]. McCullagh [10, Section 3] analyzes scalar permanent approximations for matrices with two blocks and shrinking between-block entries. We use the Gaussian law tilted by the permanent to represent the normalized raw cofactor matrix. Its phase covariance gives the limiting cofactor spectrum when the two blocks or more blocks approach decoupling.

Let $G = (\mathbb{Z}/2\mathbb{Z})^k$ be fixed, of size m, and assume n is a multiple of m. Index the rows by equal fibers over $x \in G$. Each character χ gives the sign vector s_χ on these rows and $s_0 = \mathbf{1}$ for the trivial character. We study

$$B_n = a_n I + \frac{1}{n}\sum_{\chi} \beta_{\chi,n} s_\chi s_\chi^T, \qquad a_n = 1 - \frac{1}{n}\sum_{\chi}\beta_{\chi,n}, \qquad \beta_{\chi,n} \geq 0.$$

The strengths are uniformly bounded, so $a_n > 0$ for all sufficiently large n. Character gauges are conjugations by diagonal matrices with a character's signs on the fibers. They permute the strength labels and preserve Λ.

Take independent standard circular complex Gaussians G_χ with $\mathbb{E}|G_\chi|^2 = 1$. Set

$$z_x = n^{-1/2}\sum_{\chi}\sqrt{\beta_{\chi,n}}\chi(x)G_\chi, \qquad d_x = a_n + |z_x|^2.$$

The row variables $X_i = \sqrt{a_n}\epsilon_i + z_{x_i}$, with independent standard circular ϵ_i, have covariance B_n. Wick's formula and conditional integration give $\operatorname{per} B_n = \mathbb{E}\prod_x d_x^{n/m}$. Write $\mathbb{E}_{\mathrm{t}}$ for expectation after tilting the latent Gaussian law by this product. For the raw cofactor the unpaired variables at $i \neq j$ are $\overline{X}_i X_j$ and its diagonal expectation contains $1/d_{x_i}$. It follows that

$$\frac{C(B_n)}{\operatorname{per} B_n} = \mathbb{E}_{\mathrm{t}}\big[\operatorname{diag}(a_n/d_{x_i}^2) + vv^*\big], \qquad v_i = \overline{z}_{x_i}/d_{x_i}. \tag{7}$$

Indeed the rank-one term supplies $\bar{z}_{x_i} z_{x_j}/(d_{x_i} d_{x_j})$ off the diagonal, and the diagonal completes it to $1/d_{x_i}$.

Group translations and within-fiber permutations commute with the cofactor. The fiber-constant space is diagonalized by the characters. Its orthogonal complement sees a diagonal operator of norm $O(1)$ by (7). Write $\langle f \rangle_G = m^{-1} \sum_x f_x$ for a fiber average. The uniform-mode eigenvalue has the exact defect identity

$$\mathbb{E}_{\mathrm{t}} \left\langle a_n / d_x^2 \right\rangle_G + \frac{n}{4a_n} - n\mathbb{E}_{\mathrm{t}} \left\langle \frac{(|z_x|^2 - a_n)^2}{4a_n d_x^2} \right\rangle_G - n\mathbb{E}_{\mathrm{t}} \left\langle |v_x - \langle v \rangle_G|^2 \right\rangle_G . \tag{8}$$

This follows by subtracting the fiber variance from $\left\langle |v|^2 \right\rangle_G$ and using

$$\frac{u}{(a+u)^2} = \frac{1}{4a} - \frac{(u-a)^2}{4a(a+u)^2}.$$

8 A dominant character near strength two

Theorem 3. *Suppose $\beta_{0,n} = 2 + u/\sqrt{n} + o(n^{-1/2})$ and $\beta_{\chi,n} \to \kappa_\chi < 2$ for all $\chi \neq 0$. Then*

$$\Lambda(B_n) = \frac{n}{4} + \frac{1}{2} - \frac{u^2}{16} - \sum_{\chi \neq 0} \frac{\kappa_\chi^2}{8(2-\kappa_\chi)} + o(1).$$

Proof. Rotate by the phase of G_0, let $T = |G_0|^2$, and use the independent rotated secondary Gaussians $H_\chi = X_\chi + iY_\chi$. Write $\beta = \beta_{0,n}, \kappa_{\chi,n} = \beta_{\chi,n}$ and

$$g_0 = \sqrt{\beta T/n}, \quad h_x = \sum_{\chi \neq 0} \sqrt{\kappa_{\chi,n}} \chi(x) H_\chi, \quad z_x = g_0 + n^{-1/2} h_x, \qquad X = (T - n(1 - 1/\beta))/\sqrt{n}.$$

The common rotation leaves all observables of (7) unchanged. The tilted variables X, X_χ, Y_χ converge with polynomial moments to independent centered Gaussians of variances $1, 1/2, 1/(2-\kappa_\chi)$, respectively. For domination, AM–GM and orthogonality bound the tilted product by $(a_n + [\beta T + \sum \kappa_{\chi,n} |H_\chi|^2]/n)^n$. With $\widetilde{T} = T + \sum (\kappa_{\chi,n}/\beta) |H_\chi|^2$, the Gaussian densities leave a radial factor $e^{-\widetilde{T}} (a_n + \beta \widetilde{T}/n)^n$ and secondary factors $e^{-(1-\kappa_{\chi,n}/\beta)|H_\chi|^2}$. These rates have a uniform positive lower bound. The radial factor is a translated Gamma density of shape $n+1$, with lower cutoff at $n a_n/\beta$ before translation and with center differing by $O(1)$ from $n(1-1/\beta)$ in $\widetilde{T}$. Its rescaled centered moments are bounded by the Gamma moment generating function. The shift from $\widetilde{T}$ to T is controlled by the secondary Gaussian factors. This envelope is an upper bound in integral of the same order as the true normalizer, since the local expansion below gives a comparable lower mass on a fixed compact fluctuation box.

With $q = a_n + \beta T/n = \beta + \beta X/\sqrt{n} + O(n^{-1})$, the radial log density has limit $-X^2/2$, up to its constant. In the fiber log expansion use $d_x = q + 2g_0 \operatorname{Re} h_x/\sqrt{n} + |h_x|^2/n$. Orthogonality cancels the linear term in $\operatorname{Re} h_x$. The two order-one terms give

$$\tfrac{1}{2} \sum_\chi \kappa_\chi (X_\chi^2 + Y_\chi^2) - \tfrac{1}{2} \sum_\chi \kappa_\chi X_\chi^2.$$

With the original Gaussian costs this is precisely the asserted covariance. The expansions are uniform on compact fluctuation boxes, and domination as above gives tightness and uniform integrability of polynomial moments.

On the subspace orthogonal to the uniform vector, the rank expectation in (7) has trace $n\mathbb{E}_{\mathrm{t}}\left\langle |v_x - \langle v\rangle_G|^2\right\rangle_G = O(1)$. This follows from the uniformly Lipschitz map $z \mapsto \bar{z}/(a_n + |z|^2)$ and the bounded moments of h_x. The uniform-mode eigenvalue is asymptotic to $n/4$ and is therefore the top eigenvalue for large n.

In (8) the average diagonal tends to $1/4$ and $n/(4a_n) = n/4 + (2 + \sum \kappa_\chi)/4 + o(1)$. The first deficit tends to $(u^2+4+2\sum \kappa_\chi)/16$, since $\sqrt{n}(|z_x|^2 - a_n)$ tends to $u + 2X + 2\operatorname{Re} h_x$ in moments and $d_x \to 2$. For the second deficit, the rescaled centered v_x tends to $-i \operatorname{Im} h_x/2$. This is the derivative at $z = 1, a = 1$, which vanishes in the real direction. Orthogonality and the limiting imaginary variances give the second deficit $\sum \kappa_\chi/[4(2-\kappa_\chi)]$. The positive denominators and moment bounds justify taking expectations in these expressions. Substituting the two deficits gives the formula. □

9 Global rigidity and exact maximization

Theorem 4. *Fix the Walsh group and any cap $B > 2$. For all sufficiently large n divisible by m, the maximum of $\Lambda(B_n)$ over $0 \le \beta_{\chi,n} \le B$ equals the scalar equicorrelation maximum. Every maximizer, after a character gauge, has only the trivial strength nonzero and is the scalar maximizer. In particular, the maximum is $(n+1)^2/(4n) + O(ne^{-cn})$ for some $c > 0$.*

9.1 Localization of two-term near-maximizers

Suppose first that the strengths are uniformly bounded and $\Lambda(B_n) \ge n/4 + 1/2 - o(1)$. The top eigenvector must lie in the fiber-constant space, since its complement has bounded norm. Gauge a maximizing character to the trivial one. Then (8) equals $\Lambda(B_n)$. Its two nonnegative deficits are bounded, since $a_n = 1 - O(n^{-1})$ and the diagonal term is bounded. Hence in tilted probability, for each of the fixed number of fibers,

$$|z_x|^2 = 1 + O_{\mathbb{P}}(n^{-1/2}), \qquad z_x - \langle z\rangle_G = O_{\mathbb{P}}(n^{-1/2}). \tag{9}$$

Indeed the first defect controls $(|z_x|^2 - a_n)/d_x$ at this scale, which also controls $|z_x|^2 - a_n$ by solving for it. The second defect controls $v_x - \langle v\rangle_G$; then use $z_x = d_x \bar{v}_x$ and $d_x = 2 + O_{\mathbb{P}}(n^{-1/2})$.

In rotated latent coordinates as in Theorem 3, (9) implies $\beta T/n = 1 + O_{\mathbb{P}}(n^{-1/2})$ and $\sqrt{\kappa_{\chi,n}} H_\chi = O_{\mathbb{P}}(1)$. The scaled latent vector $(G_\chi/\sqrt{n})$ is tight for bounded strengths. With $R = n^{-1}\sum |G_\chi|^2$, the Gaussian radial cost and AM–GM give the exponential tail bound with exponent $n[-R + \log(1 + BR)]$, up to a fixed-degree radial volume factor. The permanent normalizer is bounded below by $a_n^n \ge c > 0$. In particular β cannot tend to zero along near-maximizers.

Pass to a subsequence on which the strengths converge, with trivial limit $\beta_* > 0$. The scaled latent vector concentrates on the maximizers of

$$-\sum_\chi |w_\chi|^2 + \frac{1}{m}\sum_x \log\left(1 + \left|\sum_\chi \sqrt{\beta_{\chi,*}}\chi(x) w_\chi\right|^2\right).$$

This follows by uniform convergence of the exponents on compacts, the displayed tail bound, and integration of a small ball near a maximum for a lower bound. A limiting latent vector must give all z_x equal with modulus one by (9). Its positive-strength nontrivial coordinates vanish by character orthogonality, and the zero-strength coordinates vanish at a maximum since they only enter the negative quadratic term. At the trivial-axis maximizer $|w_0|^2 = 1/\beta_*$, so the radial derivative is $-1 + \beta_*/2$. Thus $\beta_* = 2$. No secondary limit can exceed two, since moving to that pure character at the same radius would then increase the exponent.

The secondary H_χ are tight when their strengths have positive limits, by (9). To see tightness when $\kappa_{\chi,n} \to 0$, condition on T and the other rotated variables. With arbitrarily high probability the trivial latent value $g_0 = \sqrt{\beta T/n}$ is in a fixed compact neighborhood of one, and all other $\sqrt{\kappa_{j,n}} H_j$ are bounded by a fixed M. Denote the field with χ omitted by $z_x^0 = g_0 + O(M/\sqrt{n})$, and write $d_x^0 = a_n + |z_x^0|^2$. The change in the log tilt is bounded above using $\log d_x - \log d_x^0 \le (d_x - d_x^0)/d_x^0$. Its sum with multiplicities is at most $C_M \sqrt{\kappa_{\chi,n}} |H_\chi| + C\kappa_{\chi,n} |H_\chi|^2$, because the linear term in the constant value of $\bar{z}_x^0/d_x^0$ cancels by the zero character sum. On $|H_\chi| \le 1$, Taylor expansion and the same cancellation give an absolute bound C_M. Thus the conditional normalizer has a positive lower bound after dividing by the value at $H_\chi = 0$, and the original Gaussian yields a common conditional Gaussian tail. This proves full secondary tightness.

Set $X_n = (T - n/\beta)/\sqrt{n}$, tight by (9). On any compact of its coordinates and the H_χ, the log of the unnormalized density divided by its central value has expansion

$$L_n X_n - \tfrac{1}{2} X_n^2 - \sum_\chi X_\chi^2 - \sum_\chi (1 - \kappa_\chi/2) Y_\chi^2 + o(1), \qquad L_n = \sqrt{n} \left(\frac{\beta}{a_n + 1} - 1 \right),$$

where κ_χ denotes the limiting secondary strengths on the subsequence. This is the same character cancellation as in Theorem 3, with radial center at n/β. If L_n were unbounded, moving any compact box a fixed distance in X_n in its sign direction would make the density on the shifted box arbitrarily larger than on the original, along a subsequence. This would force the original compact mass to vanish, contradicting tightness. Hence $\sqrt{n}(\beta - 2)$ is bounded. Pass further so that L_n converges. If some $\kappa_\chi = 2$, the local limiting unnormalized density is constant in Y_χ, so its full integral is infinite. Integrating over arbitrarily large compact boxes then makes the full normalizer divided by the central density diverge, forcing every fixed compact to have vanishing probability. This again contradicts tightness. Thus all $\kappa_\chi < 2$.

Theorem 3 now applies at every subsequential limit u of $\sqrt{n}(\beta - 2)$. Its nonnegative defect $u^2/16 + \sum_\chi \kappa_\chi^2/[8(2 - \kappa_\chi)]$ must be zero at a near-maximizer. We conclude $\sqrt{n}(\beta - 2) \to 0$ and $\max_{\chi \ne 0} \kappa_{\chi,n} \to 0$ after the chosen gauges.

9.2 A uniform strength neighborhood

Let $\beta_n^* = n b_n$, with b_n the exact scalar maximizer from Theorem 1, so $\beta_n^* = 2 + O(n^{-1})$. Parametrize the neighborhood by $\beta = \beta_n^* + u/\sqrt{n}$ and $\kappa_{\chi,n} = \psi_\chi^2$, with u, ψ in fixed sufficiently small symmetric intervals. Let $F_n(u, \psi)$ be (8) minus $n/4$. On this rectangle the uniform mode is the top eigenvalue for large n, uniformly, by the moment and complement bounds of Theorem 3. We prove

$$F_n \;\longrightarrow\; \tfrac{1}{2} - \frac{u^2}{16} - \sum_\chi \frac{\psi_\chi^4}{8(2 - \psi_\chi^2)} \quad \text{in } C^4(u, \psi).$$

Take the fixed radial center $T_0 = n(1 - 1/\beta_n^*)$, and $X = (T - T_0)/\sqrt{n}$. Put $q = a_n + \beta T/n$, $q_0 = a_n + \beta T_0/n$, and write $d_x = q + r_x$, where

$$r_x = 2 g_0 \operatorname{Re} h_x/\sqrt{n} + |h_x|^2/n, \qquad h_x = \sum_\chi \psi_\chi \chi(x) H_\chi.$$

Divide the unnormalized density in (X, H) by its value at the center. Its log tilt terms decompose as

$$n \log(q/q_0) + n \langle r_x \rangle_G / q + n \langle \log(q + r_x) - \log q - r_x/q \rangle_G .$$

The first term has polynomially bounded derivatives through order four, uniformly on the rectangle. In fact $q/q_0 = 1 + (\beta/q_0)X/\sqrt{n}$ is bounded below and each positive-order parameter derivative of β/q_0 is $O(n^{-1/2})$. The second term uses $\langle r_x \rangle_G = n^{-1} \sum \psi_\chi^2 |H_\chi|^2$. In the third term use the integral Taylor remainder with numerator r_x^2 and denominators $(q + t r_x)^2$, $0 \le t \le 1$, all bounded below by a_n^2. The factors $\sqrt{n}\, r_x$ and their parameter derivatives are polynomially bounded in (X, H), since $T/n \le C(1 + |X|)$ and derivatives of $g_0 = \sqrt{\beta T/n}$ differentiate only $\sqrt{\beta}$. This proves polynomial bounds for all differentiated density factors after taking out the density.

The diagonal contribution and the first defect in (8) have the same bounds after differentiation, since $d_x \ge 1/2$ and $\sqrt{n}(|z_x|^2 - a_n)$ is polynomially bounded with its derivatives. For the last defect, subtract the common value $v(g_0)$ before centering over the fibers. The expression $\sqrt{n}[v(z_x) - v(g_0)]$ is an integral of the real differential of v applied to h_x along the line from g_0 to z_x. The rational map $v(z) = \bar{z}/(a_n + |z|^2)$ and its fixed-order derivatives have polynomial bounds for $a_n \ge 1/2$, which also bounds all differentiated versions of this defect.

The preceding Gamma envelope is uniform on the rectangle. Its integral relative to the density value at the center is bounded in $dX\, dH$ by Stirling's formula, since the translated Gamma center is within $O(\sqrt{n})$ of T_0 and the T-to-X Jacobian is $\sqrt{n}$. Its bounded moments of all orders give uniform integrability of the polynomial derivative bounds. The true relative integral has a uniform positive lower bound on a compact fluctuation box by the local expansion. Differentiation of the normalized integrals is therefore justified, and convergence of the differentiated expansions on compact fluctuation boxes gives uniform convergence through order four. In these fixed-center coordinates, the radial limiting log tilt is $uX/4 - X^2/2$. After recentering X by $u/4$, the dominant-character calculation gives exactly the displayed limit for F_n.

The functions are separately even in every ψ_χ, by changing the sign of the corresponding H_χ. Thus C^4 convergence gives C^2 convergence on the closed strength rectangle $\kappa_\chi = \psi_\chi^2 \ge 0$. For a one-variable even function f, the first derivative of $f(\sqrt{\kappa})$ is $\frac{1}{2}\int_0^1 f''(t\sqrt{\kappa})dt$, and the second is $\frac{1}{4}\int_0^1 \int_0^1 t^2 f^{(4)}(st\sqrt{\kappa})ds\, dt$, both continuous at zero. Apply the same formula in each variable for mixed derivatives. The limiting Hessian in the strength coordinates is uniformly negative definite on a small fixed rectangle, so the exact objective is strictly concave there for large n.

Its gradient at $u = 0, \kappa = 0$ is zero exactly. The scalar coordinate is stationary by the interior scalar optimum. At the same equicorrelation point, simplicity and permutation invariance make the derivative on real symmetric zero-diagonal perturbations proportional to the sum of off-diagonal entries. Since the scalar perturbation has nonzero such sum and derivative zero, the derivative in every secondary direction is zero. Thus the unique strength maximizer in this neighborhood is the scalar point.

The capped family is compact and contains the scalar point for large n. Every maximizing sequence is two-term near-optimal; the localization above places it, after a gauge, inside the strength rectangle for sufficiently large n. The strict concavity then proves Theorem 4, including equality. The scalar value and its error follow from Theorem 1.

10 The critical quartic regime at strength one

The scale in this section concerns growth of the ratio, before the optimization near strength two. The scalar radial transition is related to incomplete Gamma asymptotics [11]. Quartic critical laws for vector spin systems [12] are another precedent for the limiting form of the exponent. Here the law controls the largest eigenvalue of a permanental cofactor.

Theorem 5. *If all $\beta_{\chi,n} \to b_\chi < 1$, then*

$$\Lambda(B_n) \longrightarrow \max_\chi (1 - b_\chi)^{-1}.$$

Suppose a nonempty character set J has $\beta_{\chi,n} = 1 + u_\chi/\sqrt{n} + o(n^{-1/2})$, with all other limits $b_\chi < 1$. Let $Y = (Y_\chi)_{\chi \in J}$ be the complex vector of density proportional to

$$\exp\left[\sum_{\chi \in J} u_\chi |Y_\chi|^2 - \frac{1}{2m} \sum_{x \in G} \left| \sum_{\chi \in J} \chi(x) Y_\chi \right|^4 \right].$$

Then

$$n^{-1/2} \Lambda(B_n) \longrightarrow \lambda_{\max}(\mathbb{E}[YY^*]) = \max_{\chi \in J} \mathbb{E}|Y_\chi|^2.$$

For one critical character of detuning u, the value is $u + \varphi(u)/\Phi(u)$, with φ, Φ the standard real normal density and distribution function.

Proof. The exact fiber identity

$$\frac{n}{m} \sum_x (d_x - 1) = \sum_\chi \beta_{\chi,n} (|G_\chi|^2 - 1)$$

follows from orthogonality. In the subcritical regime, the log tilt converges to $\sum b_\chi (|G_\chi|^2 - 1)$ at fixed G_χ. The bound $\log d_x \le d_x - 1$ gives a common Gaussian majorant with all moments, since the strengths are uniformly below one. The latent limits then have independent circular variances $(1 - b_\chi)^{-1}$. In (7) the diagonal tends to the identity. In the orthonormal fiber basis the rank expectation has matrix $(n/m)\mathbb{E}_{\mathrm{t}}[vv^*]$ using the m fiber values. Here $v_x = \bar{z}_x + o(n^{-1/2})$ in tilted L^2, so this rank contribution has limiting eigenvalues $b_\chi/(1 - b_\chi)$ by character orthogonality. The subcritical formula follows.

For critical characters set $G_\chi = n^{1/4} Y_\chi$ and leave the other Gaussians unscaled. At fixed rescaled variables, with $\psi_x = \sum_{\chi \in J} \chi(x) Y_\chi$, one has $n^{1/4} z_x \to \psi_x$. The exact linear sum cancels the order-$\sqrt{n}$ Gaussian weight of the critical variables and gives the detuning terms. The second log term tends to the quartic term in the statement. The unscaled modes have independent limiting Gaussian weights $e^{-(1-b_\chi)|G_\chi|^2}$.

For domination put $T = n^{-1/2} \sum_J \beta_{\chi,n} |Y_\chi|^2 + n^{-1} \sum_{\chi \notin J} \beta_{\chi,n} |G_\chi|^2$. AM–GM bounds the tilt by $(1 + T)^n$. Use $\log(1 + T) \le T - c \min(T^2, T)$ for $T \ge 0$. After the original Gaussian costs, the joint exponent has the upper bound

$$C\|Y\|^2 - c' \sum_{\chi \notin J} |G_\chi|^2 - c'' \min(\|Y\|^4, \sqrt{n}\, \|Y\|^2).$$

This gives a common integrable Gaussian upper envelope for all sufficiently large n, also for any fixed polynomial moment. The local exponents converge uniformly on compact sets and the limit has positive normalizing mass there. The quartic term is coercive, since $m^{-1} \sum_x |\psi_x|^4 \ge \|Y\|^4$ by Jensen and orthogonality. Thus this is the joint limiting law.

After division by $\sqrt{n}$ the bounded diagonal term of (7) vanishes. Since $n^{1/4} v_x \to \bar{\psi}_x$ in tilted L^2, the normalized rank part on the fiber basis tends to $m^{-1}\mathbb{E}[\bar{\psi}\, \bar{\psi}^*]$. Its nonzero eigenvalues are those of the conjugate of $\mathbb{E}[YY^*]$ by the orthonormal Walsh columns. The covariance is diagonal in the character basis, by invariance under group translation of the fiber index in the quartic sum. For one critical character its squared modulus has density proportional to $e^{ut - t^2/2}$ on $t \ge 0$, giving the stated truncated normal mean. □

11 Degenerate supercritical strengths

Theorem 6. *Suppose all m strengths satisfy $\beta_{\chi,n} = \gamma + \tau_\chi/n + o(n^{-1})$ with $\gamma > 1$. Put $W_{x,\chi} = \chi(x)/\sqrt{m}$ and $Q = W \operatorname{diag}(\tau_\chi) W^*$. Let $\Xi = (e^{i\theta_x})_{x\in G}$ have density proportional to*

$$\exp\left[\frac{\gamma-1}{m\gamma^2}\Xi^* Q \Xi\right]$$

with respect to uniform product angles on the torus. Then

$$\frac{\Lambda(B_n)}{n} \longrightarrow \frac{\gamma-1}{m\gamma^2}\lambda_{\max}(\mathbb{E}[\Xi\Xi^*]).$$

If $m > 1$ and all detunings are finite, this value is strictly below $(\gamma-1)/\gamma^2$. At equal detunings it is $(\gamma-1)/(m\gamma^2)$. If $m = 2$, it is

$$\frac{\gamma-1}{2\gamma^2}\left(1 + \left|\frac{I_1(\mu)}{I_0(\mu)}\right|\right), \qquad \mu = \frac{\gamma-1}{2\gamma^2}(\tau_0 - \tau_1),$$

where I_0, I_1 are modified Bessel functions.

Proof. In the latent fiber coordinates $z = (z_x)$ the covariance is $(m/n)W \operatorname{diag}(\beta_{\chi,n}) W^*$. Thus the tilted density is proportional to

$$\exp[-(n/m) z^* W \operatorname{diag}(\beta_{\chi,n}^{-1}) W^* z] \prod_x (a_n + |z_x|^2)^{n/m}.$$

At reference strength γ, the squared radii have independent translated Gamma densities with rate $n/(m\gamma)$, restricted to positive $t_x = |z_x|^2$. Their centers tend to $\gamma - 1 > 0$. The angles are uniform and independent. Expanding the inverse strengths multiplies this density by $\exp[z^* Q z/(m\gamma^2) + o(1)\|z\|^2]$. The radial reference variables have uniformly bounded exponential moments at each fixed order, from the Gamma formula and the separated lower cutoff. This controls the tilt, and radial concentration gives the displayed limiting phase density, with convergence of the moments used in the cofactor identity. Phase models of this form are studied in [13].

On the fiber basis the rank term in (7) is $(n/m)\mathbb{E}_{\mathrm{t}}[vv^*]$. Its n^{-1} limit has the spectrum asserted because v_x tends to $\sqrt{\gamma-1}\,\overline{\Xi_x}/\gamma$ in tilted expectations. The diagonal part is bounded. The phase covariance has trace m and is positive definite under the everywhere positive torus density. Thus for $m > 1$ its top eigenvalue is strictly below m. Equal detunings give a scalar Q and independent uniform angles. For two fibers the difference angle has density proportional to $e^{\mu\cos\theta}$, giving the stated I_1/I_0 Fourier moment. □

Theorems 4–6 concern different strength regimes in the same fixed Walsh family. In particular Theorem 4 determines the third term only for a fixed bounded-rank correction with balanced Walsh fibers. The unrestricted cofactor extremum remains open.

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References

[1] R. Bhatia, P. Grover, and T. Jain, *Derivatives of tensor powers and their norms*, Electron. J. Linear Algebra **26** (2013), 604–619. doi:10.13001/1081-3810.1675.

[2] R. Grone, C. R. Johnson, E. M. Sá, and H. Wolkowicz, *A note on maximizing the permanent of a positive definite Hermitian matrix, given the eigenvalues*, Linear Multilinear Algebra **19** (1986), 389–393. Author copy.

[3] Y. Lavi, *A conjecture on the permanent of positive semidefinite matrices*, 2026. Preprint, version 1.

[4] E. H. Lieb, *Proofs of some conjectures on permanents*, J. Math. Mech. **16** (1966), 127–134.

[5] M. Marcus and R. Merris, *A relation between the permanental and determinantal adjoints*, J. Austral. Math. Soc. **15** (1973), 270–271. doi:10.1017/S1446788700013173.

[6] L. Pioge, K. K. Pietrasz, B. Seron, L. Novo, and N. J. Cerf, *A logical implication between two conjectures on matrix permanents*, Linear Algebra Appl. **725** (2025), 309–318. arXiv:2508.00111.

[7] S. Schubr, *Asymptotically optimal bounds for permanental cofactor matrices*, Luminary Research Archive, ed. 4, 2026. Archive edition 4.

[8] C. Fassino, G. Pistone, and M. P. Rogantin, *Computing the moments of the complex Gaussian: Full and sparse covariance matrix*, Mathematics **7** (2019), 263. doi:10.3390/math7030263.

[9] S. Rahimi-Keshari, A. P. Lund, and T. C. Ralph, *What can quantum optics say about computational complexity theory?*, Phys. Rev. Lett. **114** (2015), 060501. arXiv:1408.3712.

[10] P. McCullagh, *An asymptotic approximation for the permanent of a doubly stochastic matrix*, Author manuscript.

[11] N. M. Temme, *Uniform asymptotic expansions of the incomplete Gamma functions and the incomplete Beta function*, Math. Comp. **29** (1975), 1109–1114. CWI archive.

[12] K. Kirkpatrick and T. Nawaz, *Asymptotics of mean-field $O(N)$ models*, arXiv:1602.03257.

[13] N. Klein, J. Orellana, S. Brincat, E. K. Miller, and R. E. Kass, *Torus graphs for multivariate phase coupling analysis*, arXiv:1910.11044.